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**SIMULATING NANOSCALE PROCESSES IN SOLIDS USING DFT AND THE
QUASICONTINUUM METHOD**

Dan Negrut*

Mathematics and Computer Science Division
Argonne National Laboratory
Argonne, Illinois 60439
Email: negrut@mcs.anl.gov

Mihai Anitescu

Mathematics and Computer Science Division
Argonne National Laboratory
Argonne, Illinois 60439
Email: anitescu@mcs.anl.gov

Todd Munson

Mathematics and Computer Science Division
Argonne National Laboratory
Argonne, Illinois 60439
Email: tmunson@mcs.anl.gov

Peter Zapol

Materials Science and Chemistry and Divisions
Argonne National Laboratory
Argonne, Illinois 60439
Email: zapol@mcs.anl.gov

ABSTRACT

This work proposes a theoretical framework for the investigation of chemical and mechanical properties of nanostructures. The methodology is based on a two-step approach to compute the electronic density distribution in and around a nanostructure and then the displacement of its nuclei. The Electronic Problem embeds interpolation and coupled cross-domain optimization techniques through a process called electronic reconstruction. In the second stage of the solution, the Ionic Problem deals with repositioning the nuclei of the nanostructure given the electronic density in the domain. It is shown that the new ionic configuration is the solution of a nonlinear system based on a first-order optimality condition when minimizing the total energy associated with the nanostructure. The long-term goal of this work is a substantial increase in the dimension of the nanostructures that can be simulated by using approaches that include accurate DFT computation. The increase in nanostructure size stems from the fact that during the solution of the Electronic Problem expensive DFT calculations are limited to a small number of subdomains; the electronic density is then reconstructed elsewhere. For the Ionic Problem, computational gains result from approximating

the dislocation of the nuclei in terms of a reduced number of representative nuclei following the quasicontinuum paradigm.

PARADIGM OF THE PROPOSED APPROACH

Nanostructures have dimensions in the range of $1 \sim 100$ nm and typically contain $10^2 \sim 10^8$ atoms. Applying the well-established Kohn-Sham DFT method [1] for nonperiodic structures of 60 atoms has led to simulations that can take up to three months to complete. When long range interactions are ignored and pseudo-potentials are used, *ab-initio* simulations have been carried out for nonmetallic structures with up to 1,500 atoms [2]. The approach that enabled the increase in the number of atoms belongs to the family of so-called $O(N)$ methods [3], which scale as N with the dimension of the problem (in this case the number of electrons).

This work is not concerned with fundamental electronic structure computation methods. Acknowledging the small-dimension constraint placed on the problem by the existing Density Functional Theory (DFT)-based methods, the goal of the proposed work is to use techniques that, by closing the spatial scale gap, render electronic structure information at the

*Address all correspondence to this author.

nanoscale. This electronic structure information is then used to investigate the chemical and mechanical properties of the material.

In the context of mechanical analysis of nanostructures, the methodology proposed follows in the steps of the quasi-continuum work proposed in [4–6]. Specifically, this is an extension of the work in [5, 6], because rather than considering a potential-based interatomic interaction that has a limited range of validity and is difficult to generalize to inhomogeneous materials, the methodology proposed uses *ab-initio* methods to provide for the particle interaction. At the same time it is a generalization of the method proposed in [4] because rather than considering each mesh discretization element to be part of a periodic and uniformly deformed infinite crystal, the proposed method treats in a generic optimization framework any structure (nonperiodic and inhomogeneous) once the electronic density distribution is available.

The electronic structure computation is approached herein as the solution of a constrained minimization problem [7]

$$\min_{\rho} E[\rho, \rho_A] \quad (1a)$$

$$\int \rho(\mathbf{r}) d\mathbf{r} = N_e \quad (1b)$$

where N_e represents the number of electrons present in the system. The solution to this problem depends parametrically on the nuclear density ρ_A , $\rho = \rho(\rho_A)$, a consequence of the Born-Oppenheimer assumption. Subsequently, the computation of the ground state of the entire system as the solution of the optimization problem

$$\min_{\rho_A} E[\rho(\rho_A), \rho_A] \quad (2)$$

From a geometric perspective two assumptions are made in order to close the gap between the subatomic-level representation of the electron density, and the nanoscale scale associated with the structures investigated: (a) there is a near regularity in the atomic compositions of the material, and (b) almost everywhere in the nanostructure the solution to the *Ionic Problem* results in only small deformations. The assumption (a) is referred to as *near-periodicity* and is the vehicle that carries first-principles computation results from micro to macro scale. This work does not build on the periodicity assumption, it merely assumes that the material displays close to periodic structure. As explained later, the *near-periodicity* assumption enables the use of interpolation for electronic structure reconstruction.

In regards to the second assumption; i.e., *small deformation*

tions, in order to formally quantify this concept, the nanostructure is considered to occupy or be contained inside an initial reference configuration $D^0 \subset \mathbb{R}^3$. The structure undergoes a change of shape described by a deformation mapping $\Phi(\mathbf{r}^0, t) \in \mathbb{R}^3$. This deformation mapping gives the location \mathbf{r} in the global Cartesian reference frame of each point \mathbf{r}^0 represented in the undeformed material frame. As indicated, the mapping might depend on time t . The variable t does not necessarily represent the time contemporary with the structure under consideration. In fact, in a quasi-state simulation framework, this variable might be an iteration index of an optimization algorithm that solves Eq.(2) in the case ρ_A is made of nuclear point charges.

The components of the deformation gradient are introduced as

$$F_{iJ} = \frac{\partial \Phi_i}{\partial r_J^0} \quad (3)$$

where upper-case indices refer to the material frame, and lower-case indices to the Cartesian global frame. Thus, $\mathbf{F} = \nabla_0 \Phi$, where ∇_0 represents the material gradient operator, and therefore the deformation of an infinitesimal material neighborhood $d\mathbf{r}^0$ about a point \mathbf{r}^0 of D^0 is expressed as

$$d\mathbf{r}_i = F_{iJ} dr_J^0 \quad (4)$$

The concept of small distortion is equivalent to requiring that the spectral radius of \mathbf{F} be sufficiently small; that is,

$$\|\nabla_0 \Phi\|_2 < \mathcal{K} \quad (5)$$

is expected to hold for almost everywhere in the domain D^0 , for a suitable chosen value of \mathcal{K} .

As a consequence of the two assumptions introduced, computational savings are anticipated due to a two-tier interpolation-based approach that will reduce the dimension of the problem. First, the electronic structure will be evaluated in some domains by interpolation using adjacent regions in which a DFT-based approach is used to accurately solve the electronic structure problem; this procedure is called *electronic density reconstruction* (EDR). Second, the position of the nuclei will be expressed in terms of the positions of a reduced set of so-called representative nuclei, *repnuclei*, in an approach similar to the one proposed in [6]. The proposed approach solves only for the position of these *repnuclei*; the position of the rest of the nuclei is then obtained by interpolation.

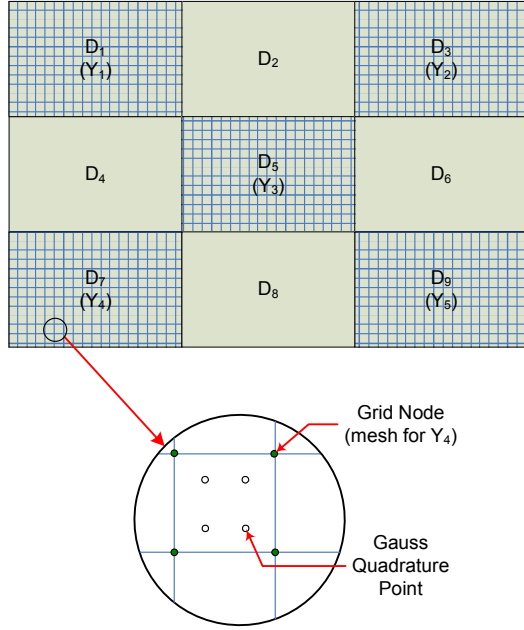


Figure 1. NANOSTRUCTURE PARTITIONED IN COMPUTATIONAL DOMAINS.

THE ELECTRONIC PROBLEM

The *Electronic Problem* refers to the computation of the electron density in a domain surrounding the nanostructure. There are two ways in which this task can be carried out: in external-mode, when third-party software is employed to this end, and internal-mode, when at least in some domain the electron density is computed directly as the solution of Eq.(1). The *Electronic Problem* is addressed by carrying out high accuracy electronic computation in a select number of domains. In the remaining domains the electron density is recovered by interpolation.

External-mode Electron Density Reconstruction

Consider, as shown in Fig.1, a two-dimensional structure $D = D_1 \cup \dots \cup D_9$ that surrounds the nanostructure. The figure presents a two-dimensional case, but the discussion applies both to the two- and the three-dimensional cases. For the external-mode EDR, accurate computation using an established code, such as NWChem [8] for instance, is used to compute ρ in the reference domains Y_1 through Y_p . The objective is to develop efficient tools that compute the solution to the electronic structure problem up to higher-order terms $O(\mathbf{F})^2 + O(\nabla_0 \mathbf{F})$ in the entire domain $D = D_1 \cup \dots \cup D_u$ that contains the nanostructure. This is equivalent to carrying out the first step of the classical homogenization technique (fluctuation reconstruction) [9].

In reference to Fig.1, in what follows, $Y = Y_1 \cup \dots \cup Y_p \subset D$; more precisely, there is an integer-to-integer mapping χ :

$\{1, \dots, p\} \rightarrow \{1, \dots, u\}$ such that $Y_j = D_{\chi(j)}$. Considering that outside D the electron density is zero, the potential generated by the total charge in the system is

$$V(\mathbf{r}) = \int_D \frac{\rho(\mathbf{r}') + \rho_A(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r}'$$

Taking for instance the reference domain Y_2 , it is important to consider separately the potential that is generated by electronic density outside this domain, whose complement is denoted by $\bar{Y}_2 = D - Y_2$, that is, for $\mathbf{r} \in D$

$$V^{ext}(\mathbf{r}; Y_2) = \int_{\bar{Y}_2} \frac{\rho(\mathbf{r}') + \rho_A(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r}' + \int_{Y_2} \frac{\rho_A(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r}' \quad (6)$$

Thus, the methodology for external-mode EDR starts by considering the external potential of Eq.(6) for each of the p reference domains Y_1 through Y_p . NWChem is applied to compute the electronic density in these domains. As far as Y_2 is concerned,

$$V(\mathbf{r}) = V^{ext}(\mathbf{r}; Y_2) + \int_{Y_2} \frac{\rho(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r}'$$

which effectively indicates that the influence of the remaining domains is perceived as the presence of an external potential in which the reference domain Y_2 is immersed. This approach however hinges upon the availability of ρ in \bar{Y}_2 . In the reference domains Y_1 through Y_p , ρ is explicitly computed, while based on the *near-periodicity* assumption, interpolation is used to recover ρ in $D - Y$. For ρ_i on domain D_i a set of weights ϑ is considered that depends exclusively on the type of interpolation considered (linear, quadratic, etc.). Thus,

$$\rho_i(\Phi(\mathbf{r}^{0'}, t)) = \sum_{\alpha=1}^p \vartheta_{\alpha}(i) \rho_{\alpha}(\Phi(\mathbf{r}^{0'} + \mathbf{T}_{i\alpha}, t)) \quad (7)$$

where the vector $\mathbf{T}_{i\alpha}$ is the translation vector that based on the periodicity assumption takes the point $\mathbf{r}^{0'}$ in domain D_i^0 to its image in the domain Y_{α} . Here, a zero superscript indicates the undeformed domain, and by convention, Greek subscripts are used to index quantities associated with a reconstruction domain Y_{α} , $\alpha = 1, \dots, p$. Note that in order to define the deformation field, an appropriate representation for $\Phi(\mathbf{r}^{0'}, t)$ is necessary. Following the quasicontinuum paradigm [5, 6], this mapping is defined

based on the relative displacement of a subset of nuclei in the nanostructure:

$$\Phi(\mathbf{r}^{0'}, t) = \sum_{A \in \mathcal{B}} \phi(\mathbf{r}^{0'} | \mathbf{R}_A^0) \Phi(\mathbf{R}_A^0, t) \quad (8)$$

The deformation needs to be represented only at the points \mathbf{R}_A^0 , $A \in \mathcal{B}$, and is then reconstructed by interpolation at the other points of the space, by using the shape functions $\phi(\cdot, \cdot)$. This aspect of the presentation is going to be covered in more detail when discussing the *Ionic Problem*.

Note that not all domains need to be included in the electronic density reconstruction. If a domain Y_γ has a defect, the electronic density is severely distorted away from the near-periodicity assumption and should not be included in the reconstruction process in neighboring domains. This translates in taking $\vartheta_\gamma(i) = 0$ for a domain $D_i \subset D - Y$ when carrying out interpolation (the density ρ_γ is nevertheless used for reconstructing the external potential $V^{ext}(\mathbf{r}; Y_\alpha)$ for each $\alpha \neq \gamma$).

Concluding on the external-mode EDR, based on an initial electron density distribution, three steps that effectively amount to a nonlinear Gauss-Jacobi method are iteratively taken to reconstruct the electron density in D : (1) for each Y_α generate the external potential as in Eq.(6); (2) carry out accurate DFT computation in each domain Y_α ; (3) use Eq.(7) to evaluate ρ in $D - Y$. This iterative process stops when the change in electron density between successive iterations becomes smaller than a threshold value.

Internal-mode Electron Density Reconstruction

The discussion for the internal-mode EDR starts with the premise that for a given ionic distribution, the electronic energy is expressed as

$$E(\rho, \rho_A) = \int \Theta^1(\rho, \rho_A, \mathbf{r}) d\mathbf{r} + \iint \Theta^2(\rho, \rho_A, \mathbf{r}; \rho, \rho_A, \mathbf{r}') d\mathbf{r} d\mathbf{r}' \quad (9)$$

This representation is commonly used in conjunction with the so-called Orbital-Free DFT (OFDFT) method [10]. Here $\Theta^{1,2}$ are the relevant energy density functionals; ρ is the electronic density; and ρ_A is the nuclear density, which may include delta functions. The first term typically includes the kinetic energy and an exchange-correlation term, whereas the second integral includes all pairwise interactions. Details regarding the definition of these terms are provided by several authors [11–13].

As an example in this paper the proposed methodology is applied in conjunction with Thomas-Fermi DFT [14, 15]. The Thomas-Fermi functional has well-known severe accuracy limitations. It provides, however, a simple vehicle that allows the

introduction of an otherwise involved reconstruction methodology. The Thomas-Fermi-based energy functional assumes the form

$$E[\rho, \{\mathbf{R}_A\}] = E_{ne}[\rho, \{\mathbf{R}_A\}] + J[\rho] + K[\rho] + T[\rho] + V_{nn}(\{\mathbf{R}_A\}) \quad (10)$$

where

$$E_{ne}[\rho, \{\mathbf{R}_A\}] = - \sum_{A=1}^M \int \frac{Z_A \rho(\mathbf{r})}{\|\mathbf{R}_A - \mathbf{r}\|} d\mathbf{r} \quad (11a)$$

$$J[\rho] = \frac{1}{2} \iint \frac{\rho(\mathbf{r}) \rho(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r} d\mathbf{r}' \quad (11b)$$

$$T[\rho] = C_F \int \rho^{\frac{5}{3}}(\mathbf{r}) d\mathbf{r} \quad (11c)$$

$$K[\rho] = -C_x \int \rho^{\frac{4}{3}}(\mathbf{r}) d\mathbf{r} \quad (11d)$$

$$V_{nn}(\{\mathbf{R}_A\}) = \sum_{A=1}^M \sum_{B=A+1}^M \frac{Z_A Z_B}{\|\mathbf{R}_A - \mathbf{R}_B\|} \quad (11e)$$

Here $C_F = \frac{3}{10}(3\pi^2)^{2/3}$, and $C_x = \frac{3}{4}\left(\frac{3}{\pi}\right)^{1/3}$, and the following notation is used:

- E_{ne} - energy corresponding to nucleus-electron interaction
- J - Coulomb energy
- K - exchange energy
- T - kinetic energy
- V_{nn} - internuclear interaction energy
- Z_A - atomic number associated with nucleus A
- \mathbf{r}_i - global position of electron i
- \mathbf{R}_A - global position of nucleus of atom A
- $\int(\cdot)$ without integration limits - an integral over the entire domain.

The expression of the energy functional of Eq.(10) justifies the notation used in Eq.(9): the kinetic, exchange, and nuclear-electronic energy are represented through the Θ^1 term; the electron-electron interaction is associated with the term Θ^2 . In the Thomas-Fermi case, the optimization problem of Eq.(1) depends parametrically on the positions of the nuclei:

$$\min_{\rho} \quad E(\rho; \{\mathbf{R}_A\}) + \lambda \left(\int \rho d\mathbf{r} - N \right) \quad (12a)$$

$$\text{s.t.} \quad \int \rho d\mathbf{r} - N = 0 \quad (12b)$$

In a domain decomposition framework, it can be shown [16]

that solving the above optimization problem in D is equivalent to solving a set of smaller optimization problems on domains D_i . For this define

$$E_i[\rho_i, \lambda_i; \bar{\rho}_i, \{R_A\}] = C_F \int_{D_i} \rho_i^{\frac{5}{3}}(\mathbf{r}) d\mathbf{r} - C_x \int_{D_i} \rho_i^{\frac{4}{3}}(\mathbf{r}) d\mathbf{r} \quad (13)$$

$$+ \int_{D_i} \int_{D-D_i} \frac{\rho_i(\mathbf{r}) \bar{\rho}_i(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r} d\mathbf{r}' + \frac{1}{2} \int_{D_i} \int_{D_i} \frac{\rho_i(\mathbf{r}) \bar{\rho}_i(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r} d\mathbf{r}'$$

$$- \sum_{A=1}^M \int_{D_i} \frac{Z_A \rho_i(\mathbf{r})}{\|\mathbf{R}_A - \mathbf{r}\|} d\mathbf{r} + \lambda_i \int_{D_i} \rho_i d\mathbf{r}$$

where $\bar{\rho}_i$ denotes the electronic density in \bar{D}_i , $i = 1, \dots, u$ used as in Eq.(6) to generate the external potential in which the domain D_i is considered immersed.

The nonlinear-system approach. The equivalent optimality conditions for the subdomain optimization problems [16] are

$$\nabla_{\rho_i} E_i(\rho_i, \lambda_i; \bar{\rho}_i, \{\mathbf{R}_A\}) = 0, \quad i = 1, \dots, u \quad (14a)$$

$$\lambda_1 = \dots = \lambda_u \quad (14b)$$

$$\int \rho d\mathbf{r} - N_e = 0 \quad (14c)$$

Following the reconstruction paradigm, this accurate computation of the electron density is only going to be carried out in the reference subdomains Y_1 through Y_p (that is, $D_{\chi(1)}$ through $D_{\chi(p)}$). The first order optimality conditions for a generic domain $Y_\alpha \in Y$ assume the form [16]

$$\frac{5}{3} C_F \rho_\alpha^{\frac{2}{3}}(\mathbf{r}) - \frac{4}{3} C_x \rho_\alpha^{\frac{1}{3}}(\mathbf{r}) + \sum_{i=1}^u \int_{D_i} \frac{\rho_i(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r}' \quad (15)$$

$$- \sum_{A=1}^M \frac{Z_A}{\|\mathbf{r} - \mathbf{R}_A\|} + \lambda = 0$$

Next, the density ρ_i on domain D_i is expressed in terms of reconstruction densities $\rho_\alpha \in Y_\alpha, \alpha \in \{1, \dots, p\}$ based on Eq.(7). Taking into account the deformation of the structure,

$$\int_{D_i} \frac{\rho_i(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} d\mathbf{r}' = \sum_{\alpha=1}^p \int_{Y_\alpha^0} \rho_\alpha(\Phi(\mathbf{r}^{0'}, t)) \tilde{K}_{i\alpha}(\mathbf{r}^0, \mathbf{r}^{0'}) d\mathbf{r}^{0'} \quad (16a)$$

$$\tilde{K}_{i\alpha}(\mathbf{r}^0, \mathbf{r}^{0'}) = \frac{\vartheta_\alpha(i) \|\mathbf{F}(\mathbf{r}^{0'} - \mathbf{T}_{i\alpha}, t)\|}{\|\Phi(\mathbf{r}^0, t) - \Phi(\mathbf{r}^{0'} - \mathbf{T}_{i\alpha}, t)\|} \quad (16b)$$

Define for $\mathbf{r}^0 \in Y_\alpha^0$

$$K_{\alpha\gamma}(\mathbf{r}^0, \mathbf{r}^{0'}) = \sum_{i=1}^u \tilde{K}_{i\gamma}(\mathbf{r}^0, \mathbf{r}^{0'}) \quad (17a)$$

Then, Eq.(15) yields

$$\frac{5}{3} C_F \rho_\alpha^{\frac{2}{3}}(\Phi(\mathbf{r}^0, t)) + \sum_{\gamma=1}^p \int_{Y_\gamma^0} K_{\alpha\gamma}(\mathbf{r}^0, \mathbf{r}^{0'}) \rho_\gamma(\Phi(\mathbf{r}^{0'}, t)) d\mathbf{r}^{0'} \quad (17b)$$

$$- \frac{4}{3} C_x \rho_\alpha^{\frac{1}{3}}(\Phi(\mathbf{r}^0, t)) - \sum_{A=1}^M \frac{Z_A}{\|\Phi(\mathbf{r}^0, t) - \Phi(\mathbf{R}_A, t)\|} + \lambda = 0$$

which should hold for any $\mathbf{r}^0 \in Y_\alpha$. Finally, since $\rho \geq 0$, a new function η is introduced such that

$$\rho(\Phi(\mathbf{r}^0, t)) = \eta^s(\mathbf{r}^0, t) \quad (18a)$$

where $s \geq 4$ is an even integer; a recommended value is $s = 4$, and this will enforce only non-negative values for the density ρ . This new function must then satisfy in the subdomain Y_α the following integral equations:

$$\frac{5}{3} C_F \eta_\alpha^{\frac{2s}{3}} - \frac{4}{3} C_x \eta_\alpha^{\frac{s}{3}} + \sum_{\gamma=1}^p \int_{Y_\gamma^0} K_{\alpha\gamma}(\mathbf{r}^0, \mathbf{r}^{0'}) \eta_\gamma^s(\mathbf{r}^{0'}, t) d\mathbf{r}^{0'} \quad (18b)$$

$$- \sum_{A=1}^M \frac{Z_A}{\|\Phi(\mathbf{r}^0, t) - \Phi(\mathbf{R}_A, t)\|} + \lambda = 0$$

The algorithm at this point calls for the solution of a nonlinear system of integral equations in $\rho_\alpha, \alpha = 1, \dots, p$. In order to solve this system, the reconstruction domains Y_α are meshed by using hexahedrons. These meshes are denoted in what follows by \mathcal{G}_1 through \mathcal{G}_p , and they are associated with Y_1 through Y_p , respectively.

The direct numerical solution of the nonlinear system of integral equations becomes intractable in Cartesian coordinates because of the singularity when the grid points in a mesh \mathcal{G}_α approach a nuclei of location \mathbf{R}_A (see Eq.(18b)). When approached in spherical coordinates in a three-dimensional representation this apparent singularity is nonexistent [16]. Below, a potential-smoothing step is introduced to address the situation when $\mathbf{r}^0 \rightarrow \mathbf{R}_A^0$. Compared to the original term $\|\Phi(\mathbf{r}^0, t) - \Phi(\mathbf{R}_A^0, t)\|^{-1}$, the δ -smoothing function

$$S_\delta(\mathbf{r}^0, \mathbf{R}_A^0, t) = \frac{1 - e^{-\frac{\|\Phi(\mathbf{r}^0, t) - \Phi(\mathbf{R}_A^0, t)\|}{\delta}}}{\|\Phi(\mathbf{r}^0, t) - \Phi(\mathbf{R}_A^0, t)\|} \quad (19)$$

behaves similarly for large values of $\|\Phi(\mathbf{r}^0, t) - \Phi(\mathbf{R}_A^0, t)\|$ and δ small but positive, but it converges to $\frac{1}{\delta}$ rather than going to infinity when $\mathbf{r}^0 \rightarrow \mathbf{R}_A^0$. Thus, the smoothing process applied to Eq.(18b) leads to

$$\begin{aligned} \frac{5}{3}C_F\eta_\alpha^{\frac{2s}{3}} - \frac{4}{3}C_x\eta_\alpha^{\frac{s}{3}} + \sum_{\gamma=1}^p \int_{Y_\gamma^0} K_{\alpha\gamma}(\mathbf{r}^0, \mathbf{r}^{0'}) \eta_\gamma^s(\mathbf{r}^{0'}, t) d\mathbf{r}^{0'} \quad (20) \\ - \sum_{A=1}^M Z_A S_\delta(\mathbf{r}^0, \mathbf{R}_A^0, t) + \lambda = 0 \end{aligned}$$

To make the presentation simpler, the following notation is introduced:

- $\eta_{\beta j}$ – the value of η at the node j of grid \mathcal{G}_β
- τ – a generic grid discretization cell of volume $\|\tau\|$
- $\mathcal{V}(\tau)$ – the set of vertices associated with cell τ (four for a tetrahedron, eight for an hexahedron, etc.)
- $|\mathcal{G}_\alpha|$ – the number of grid points in \mathcal{G}_α
- Y_γ^0 – undeformed reconstruction domain meshed with \mathcal{G}_γ ;
- $Y_\gamma = \cup_{\tau \in \mathcal{G}_\gamma} \tau$

After discretization, the integral equation above yields at an arbitrary grid node $i \in \mathcal{G}_\alpha$ of location $\mathbf{r}_i^0 \in Y_\alpha$,

$$\begin{aligned} \frac{5}{3}C_F\eta_{\alpha i}^{\frac{2s}{3}} + \sum_{\gamma=1}^p \left[\sum_{\tau \in \mathcal{G}_\gamma} \int_{\tau} K_{\alpha\gamma}(\mathbf{r}_i^0, \mathbf{r}^{0'}) \eta_\gamma^s(\mathbf{r}^{0'}, t) d\mathbf{r}^{0'} \right] \quad (21) \\ - \frac{4}{3}C_x\eta_{\alpha i}^{\frac{s}{3}} - \sum_{A=1}^M Z_A S_\delta(\mathbf{r}_i^0, \mathbf{R}_A^0, t) + \lambda = 0 \end{aligned}$$

The integral on τ is performed by q -point Gaussian numerical

quadrature with weights w_l :

$$\int_{\tau} K_{\alpha\gamma}(\mathbf{r}_i^0, \mathbf{r}^{0'}) \eta_\gamma^s(\mathbf{r}^{0'}, t) d\mathbf{r}^{0'} \approx \|\tau\| \sum_{l=1}^q w_l K_{\alpha\gamma}(\mathbf{r}_i^0, \mathbf{r}_l^{0'}) \eta_\gamma^s(\mathbf{r}_l^{0'}, t)$$

Figure 1 shows in the two-dimensional case a mesh cell and the quadrature points. As indicated in this figure, \mathbf{r}_i^0 describes the position of the grid nodes; the interior points (quadrature points) are located at $\mathbf{r}_l^{0'}$. The abscissas $\mathbf{r}_l^{0'}$ of the quadrature points are different from the mesh (grid) points, and the value of the unknown function η at these abscissas is obtained by interpolation. Interpolation at point $\mathbf{r}_l^{0'} \in \tau$, using a set of shape functions ϕ_d associated with the nodes $d \in \mathcal{V}(\tau)$, yields

$$\eta_\gamma^s(\mathbf{r}_l^{0'}, t) \approx \sum_{d \in \mathcal{V}(\tau)} \eta_{\gamma d}^s \phi_d(\mathbf{r}_l^{0'}, t) = \sum_{d \in \mathcal{V}(\tau)} \eta_{\gamma d}^s \phi_d^l$$

where ϕ_d^l are constants that can be precomputed. If one defines for $\mathbf{r}^0 \in Y_\alpha$ and $\mathbf{r}_l^{0'} \in Y_\gamma$

$$k_{\alpha\gamma d}(\mathbf{r}^0) = \sum_{l=1}^q w_l \phi_d^l K_{\alpha\gamma}(\mathbf{r}^0, \mathbf{r}_l^{0'}) , \quad (22a)$$

the discretized form of the integral equation expressed at grid node $i \in \mathcal{G}_\alpha$ of location $\mathbf{r}_i^0 \in Y_\alpha$ becomes

$$\begin{aligned} \frac{5}{3}C_F\eta_{\alpha i}^{\frac{2s}{3}} - \frac{4}{3}C_x\eta_{\alpha i}^{\frac{s}{3}} + \sum_{\gamma=1}^p \left[\sum_{\tau \in \mathcal{G}_\gamma} \|\tau\| \sum_{d \in \mathcal{V}(\tau)} k_{\alpha\gamma d}(\mathbf{r}_i^0) \eta_{\gamma d}^s \right] \quad (22b) \\ - \sum_{A=1}^M Z_A S_\delta(\mathbf{r}_i^0, \mathbf{R}_A^0, t) + \lambda = 0 \end{aligned}$$

By denoting the left side of Eq.(22b) by $P_{\alpha i}(\eta)$, where $\eta = (\eta_{11}, \eta_{12}, \dots, \eta_{p1}, \eta_{p2}, \dots)^T$, the nonlinear system of equations that should be solved becomes

$$P_{\alpha i}(\eta) = 0 \quad (23)$$

for $\alpha \in \{1, 2, \dots, p\}$, $i = 1, \dots, |\mathcal{G}_\alpha|$.

One additional equation is added to the set of $\sum_{\alpha=1}^p |\mathcal{G}_\alpha|$ equations above, and it follows from the charge constraint of Eq.(1b). The central idea is again to use the electronic density in the reconstruction domains Y_α to express the electronic density in the whole domain D . Skipping the intermediary steps, this yields

$$\int_D \rho(\mathbf{r}) d\mathbf{r} = \sum_{\alpha=1}^p \int_{Y_{\alpha}^0} \eta_{\alpha}^s(\mathbf{r}^0, t) \hat{K}_{\alpha}(\mathbf{r}^0, t) d\mathbf{r}^0 \quad (24a)$$

$$\hat{K}_{\alpha}(\mathbf{r}^0, t) = \sum_{i=1}^u \vartheta_{\alpha}(i) |\mathbf{F}(\mathbf{r}^0 - \mathbf{T}_{i\alpha}, t)| \quad (24b)$$

For the charge constraint equation, using for the evaluation of the integral on a cell of the undeformed grid Y_{α}^0 the same quadrature rule and using the same interpolation method to evaluate the function at the quadrature points yields:

$$\sum_{\alpha=1}^p \left[\sum_{\tau \in Y_{\alpha}^0} \sum_{d \in \mathcal{V}(\tau)} \eta_{\alpha d}^s \hat{k}_{\alpha d} \right] - N_e = 0 \quad (25a)$$

$$\hat{k}_{\alpha d} = \sum_{l=1}^q w_l \phi_d^l \hat{K}_{\alpha}(\mathbf{r}_l^0, t) \quad (25b)$$

If a Newton-type method is considered for the solution of the nonlinear system of Eqs.(23) and (25a), the partials are computed as

$$\frac{\partial P_{\alpha i}}{\partial \eta_{\alpha i}} = \frac{10s}{9} C_F \eta_{\alpha i}^{\frac{2s-3}{3}} - \frac{4s}{9} C_x \eta_{\alpha i}^{\frac{s-3}{3}} + \sum_{\tau \in Y_{\alpha}^0} \|\tau\| \delta_{\tau i}^{\alpha} k_{\alpha \alpha i}(\mathbf{r}_i^0) \eta_{\alpha i}^{s-1} \quad (26a)$$

where $\delta_{\tau i}^{\alpha} = 1$ if for $\tau \in \mathcal{G}_{\alpha}$, $i \in \mathcal{V}(\tau)$, and $\delta_{\tau i}^{\alpha} = 0$ otherwise. When $i \neq j$ or $\alpha \neq \beta$,

$$\frac{\partial P_{\alpha i}}{\partial \eta_{\beta j}} = \sum_{\tau \in Y_{\beta}^0} \|\tau\| \delta_{\tau j}^{\beta} k_{\alpha \beta j}(\mathbf{r}_i^0) \eta_{\beta j}^{s-1} \quad (26b)$$

Likewise,

$$\frac{\partial P_{00}}{\partial \eta_{\beta j}} = \sum_{\alpha=1}^p \sum_{\tau \in Y_{\alpha}^0} s \eta_{\alpha d}^{s-1} \delta_{\tau d}^{\alpha} \hat{k}_{\alpha d} \quad (26c)$$

where, by convention, $P_{00}(\eta)$ is a notation for the left side of Eq.(25a).

The Optimizations approach. In this approach the interpolation of Eq.(7) is used to create a reduced energy functional that depends only on the electron density ρ_{α} , in the reconstruction domains Y_{α} , $\alpha = 1, \dots, p$. This reduced energy functional is

then minimized.

For $\alpha, \gamma = 1, 2, \dots, p$, $\mathbf{r}^0 \in Y_{\alpha}^0$, $\mathbf{r}^{0'} \in Y_{\gamma}^0$, define

$$\tilde{J}_{\alpha\gamma}(\mathbf{r}^0, \mathbf{r}^{0'}) = \sum_{i=1}^u \sum_{j=1}^u v_{\alpha}(i) v_{\gamma}(j) \frac{|F(\mathbf{r}^0 - \mathbf{T}_{i\alpha}, t)| \cdot |F(\mathbf{r}^{0'} - \mathbf{T}_{j\gamma}, t)|}{\|\Phi(\mathbf{r}^0 - \mathbf{T}_{i\alpha}, t) - \Phi(\mathbf{r}^{0'} - \mathbf{T}_{j\gamma}, t)\|}$$

$$\tilde{L}_{\alpha}(\mathbf{r}^0) = \sum_{A=1}^M \sum_{i=1}^u v_{\alpha}(i) \frac{|F(\mathbf{r}^0 - \mathbf{T}_{i\alpha}, t)|}{\|\Phi(\mathbf{r}^0 - \mathbf{T}_{i\alpha}, t) - \Phi(\mathbf{R}_A^0, t)\|}$$

$$\tilde{M}_{\alpha}(\mathbf{r}^0) = \sum_{i=1}^u v_{\alpha}(i) |F(\mathbf{r}^0 - \mathbf{T}_{i\alpha}, t)|$$

Based on Eq.(7), several of the terms in Eq.(10) become functions of densities in the representative domains:

$$J(\rho) = \frac{1}{2} \sum_{\alpha=1}^p \sum_{\gamma=1}^p \int_{Y_{\alpha}^0} \int_{Y_{\gamma}^0} \tilde{J}_{\alpha\gamma}(\mathbf{r}^0, \mathbf{r}^{0'}) \rho_{\alpha}(\Phi(\mathbf{r}^0, t)) \rho_{\gamma}(\Phi(\mathbf{r}^{0'}, t)) d\mathbf{r}^0 d\mathbf{r}^{0'}$$

$$E_{ne}(\rho) = - \sum_{\alpha=1}^p \int_{Y_{\alpha}^0} \tilde{L}_{\alpha}(\mathbf{r}^0) \rho_{\alpha}(\Phi(\mathbf{r}^0, t)) d\mathbf{r}^0,$$

$$\int \rho d\mathbf{r} = \sum_{\alpha=1}^p \int_{Y_{\alpha}^0} \tilde{M}_{\alpha}(\mathbf{r}^0) \rho_{\alpha}(\Phi(\mathbf{r}^0, t)) d\mathbf{r}^0$$

The difficult part has to do with the kinetic energy and exchange terms $T[\rho]$, $K[\rho]$ whose dependence on the density is not linear and, outside the Thomas-Fermi theory, not even simple to state. Assume that the latter terms are described by a univariate density $\Theta^1(\rho, r)$, as indicated by the first term of Eq.(9). For instance, for the Thomas-Fermi representation,

$$\Theta^1(\rho, \mathbf{r}) = C_F \rho^{\frac{5}{3}}(\mathbf{r}) - C_x \rho^{\frac{4}{3}}(\mathbf{r})$$

which then leads to

$$T[\rho] + K[\rho] \approx \sum_{\alpha=1}^p \int_{Y_{\alpha}^0} \tilde{M}_{\alpha}(\mathbf{r}^0) \Theta^1(\rho_{\alpha}, \Phi(\mathbf{r}^0, t)) d\mathbf{r}^0$$

With these approximations and definitions and referring back to Eq.(10), the following electronic structure computation problem is defined:

$$\min E^{IO}(\rho), \quad \text{subject to } \int \rho = N, \quad (27)$$

where we use the superscript “IO” to denote the “interpolate-and-optimize” approach, and E^{IO} represents the quantity obtained in Eq.(10) after expressing the energy as a function of electron densities in the reconstruction subdomains only. The evaluation of the energy E^{IO} on a grid, and details of the optimization process (gradient computation) is detailed in an upcoming publication [17].

Nonlinear equations vs. optimization approaches

Recall that optimality conditions followed by interpolation eventually led to the nonlinear system of Eqs.(22b) and (24a). One can immediately prove that, in aggregate, this system does not represent the first-order conditions of an optimization problem. That issue is a bit unsettling because solving optimization problems is typically a more robust process than is solving equivalent nonlinear equations, since any local minimum of the optimization problem satisfies the nonlinear equation of its optimality conditions. When only a nonlinear system is available, a local minimum of the residual is not necessarily a solution of the nonlinear system. It is therefore important to assess whether there exists an optimization problem that is equivalent, at least up to the leading order of the homogenization error, with the nonlinear system (see [16] for a more detailed discussion on this topic).

In an abstract formulation, we have the following problem:

$$\min_{x_1, x_2} f(x_1, x_2)$$

where the variables x_1 correspond to the representative degrees of freedom, whereas x_2 correspond to the rest of the degrees of freedom, and an operator T relates x_2 to x_1 . Thus, in EDR the representative degrees of freedom are the ones used to parametrize the electron density in the representative domains Y_α , $\alpha = 1, 2, \dots, p$; the mapping $T(\cdot)$ is the interpolation-based operator from (7). Likewise, in the quasicontinuum method [6], the representative degrees of freedom are the positions of the representative nuclei *repnuclei*; the location of the remaining nuclei, abstractly denoted by x_2 , is then obtained as $x_2 = T(x_1)$, where $T(x_1)$ is the piecewise linear interpolation mapping with nodes at the *repnuclei*.

Based on this observation, one can formulate the nonlinear equation

$$\nabla_{x_1} f(x_1, x_2), x_2 = T(x_1) \quad (28)$$

which will provide the same solution as the original problem. However, the problem is an equilibrium problem with equilibrium constraints rather than a minimization problem. Furthermore, it immediately results using the chain rule that the opti-

mization problem

$$\min_{x_1} f(x_1, T(x_1)) \quad (29)$$

has the same solution as the previous two, provided that the reduced Hessian is positive definite, which should be true if the original Hessian was positive definite, and the interpolation mapping is full rank. This observation presents the advantage that one solves an optimization problem as opposed to a system of nonlinear equations, with better global convergence safeguards. When there are many local minima, this should help avoid the points that do not have the correct inertia of the Hessian.

The case in Eq.(28) corresponds to the internal nonlinear-system approach, whereas the case in Eq.(29) corresponds to the internal optimization approach. The following result settles in the positive the question of whether the two approaches are equivalent in the limit of the ansatz $x_2 = T(x_1)$ [16]:

Theorem 1. *Assume that the solution $x^* = (x_1^*, x_2^*)$ of the original optimization problem satisfies $\|x_2^* - T(x_1^*)\| \ll 1$; therefore the multiscale ansatz is not perfect but is merely very good. Then the solution \tilde{x}_1 of the nonlinear equation and \hat{x}_1 of the reduced optimization problem satisfy*

$$\|x_1^* - \tilde{x}_1\| = O(\|x_2^* - T(x_1^*)\|^2) \quad \|x_1^* - \hat{x}_1\| = O(\|x_2^* - T(x_1^*)\|^2)$$

NANOSTRUCTURE SHAPE INVESTIGATION

The optimization of the geometry of a nanostructure (called hereafter the *Ionic Problem*), to find the most stable shape reduces to minimizing the total energy given a ground state electronic energy E_e as a function of the position of the nuclei. More precisely, the equilibrium configuration of a nanostructure is provided by that distribution of the nuclei that minimizes the energy

$$E_{tot} = E_e + E_{nn} \quad (30)$$

where E_{nn} is the nucleus-nucleus interaction energy and, central to this development, E_e is the electronic ground-state energy for the considered nuclear distribution.

Following the Born-Oppenheimer assumption, the electronic energy depends parametrically on the positions of the nuclei through the dependence of the electronic density on the nuclei positions. Thus, in a general form (that has the Thomas-Fermi of Eq.(10) as a subcase),

$$E_e = T[\rho(\mathbf{r})] + E^{Har}[\rho(\mathbf{r})] + E^{xc}[\rho(\mathbf{r})] + \int \rho(\mathbf{r}) V_{ext}(\mathbf{r}; \{\mathbf{R}_A\}) d\mathbf{r} \quad (31)$$

where $T[\rho(\mathbf{r})]$ is the kinetic energy functional, $E^{Har}[\rho(\mathbf{r})]$ is the electron-electron Coulomb repulsion energy, $E^{xc}[\rho(\mathbf{r})]$ is the exchange and correlation energy, and $V_{ext}(\mathbf{r}; \{\mathbf{R}_A\})$ is the ionic potential, which parametrically depends on the distribution of the nuclei $\{\mathbf{R}_A\}$. The explicit dependence of $T[\rho(\mathbf{r})]$ and $E^{xc}[\rho(\mathbf{r})]$ on the density $\rho(\mathbf{r})$ is typically not available, and consequently it is approximated in some fashion [10, 18–20], an issue beyond the scope of this document. According to the Hohenberg-Kohn theorem [7], the electronic density is such that it minimizes E_e subject to the charge conservation constraint of Eq.(1b). The following results is proved in [16]:

Theorem 2. *Consider the optimization problem*

$$\min_{\{\mathbf{R}_A\}} E_{tot} = E_e + E_{nn} \quad (32a)$$

subject to the constraint that for a nuclear configuration $\{\mathbf{R}_A\}$ the energy E_e is the electronic ground-state energy, and the electronic density $\hat{\rho}$ that realizes this electronic ground energy additionally satisfies the charge constraint equation of Eq.(1b). Under these assumptions, the first-order optimality conditions for the optimization problem of Eq.(32a) lead to

$$\mathbf{F}_K = \frac{\partial E_{ext}}{\partial \mathbf{R}_K} + \frac{\partial E_{nn}}{\partial \mathbf{R}_K} = \mathbf{0} \quad (32b)$$

where \mathbf{F}_K is interpreted as the force acting on nucleus K , and

$$E_{ext}(\mathbf{r}; \{\mathbf{R}_A\}) = - \sum_{A=1}^M \int \hat{\rho}(\mathbf{r}) V_{ext}(\mathbf{r}; \{\mathbf{R}_A\}) d\mathbf{r} \quad (32c)$$

$$\begin{aligned} &= - \sum_{A=1}^M \int \frac{Z_A \hat{\rho}(\mathbf{r})}{|\mathbf{r} - \mathbf{R}_A|} d\mathbf{r} \\ E_{nn} &= \frac{1}{2} \sum_{A=1}^M \sum_{B=A+1}^M \frac{Z_A Z_B}{|\mathbf{R}_A - \mathbf{R}_B|} \end{aligned} \quad (32d)$$

For each nucleus K in the system, Eq.(32b) leads to the condition

$$\int \hat{\rho}(\mathbf{r}) \frac{\mathbf{r} - \mathbf{R}_K}{|\mathbf{r} - \mathbf{R}_K|^{\frac{3}{2}}} d\mathbf{r} + \sum_{A=1, A \neq K}^M Z_A \frac{\mathbf{R}_A - \mathbf{R}_K}{|\mathbf{R}_A - \mathbf{R}_K|^{\frac{3}{2}}} = \mathbf{0} \quad (33)$$

A couple of remarks are in order. First, the value of the above theorem is that it allows to solve the nuclear equilibrium problem by using only the solution of the electronic density problem, and not the values and the derivatives of the kinetic and exchange energy functionals. Therefore, an entirely nontransparent encapsulation of the electronic structure problem can be used, which allows the proposed approach to work well with legacy codes that do not provide all the needed derivatives. The key observation is that once the electronic density is available, the equilibrium conditions of Eq.(33) can be imposed right away. Whether the electronic structure computation is done with third-party software is irrelevant; moreover, there is no need to know the explicit dependence of the energy E_e on the electronic density $\rho(\mathbf{r})$. Second, as suggested in [21], the one-atom conditions of Eq.(32b) can be replaced by cluster conditions, an alternative that will be explored in the future.

When a local quasicontinuum approach is used, the condition of Eq.(33) is imposed only for *repnuclei*; that is, only for $K \in \mathcal{B}$ (see Eq.(8)). The position of the rest of the atoms in the system is then expressed in terms of the position of the *repnuclei*. The *repnuclei* become the nodes of an atomic mesh, and interpolation is used to recover the position of the remaining nuclei. For instance, if the atomic mesh is denoted by \mathcal{M} , τ is an arbitrary cell in this mesh, $\mathcal{V}(\tau)$ represents the set of the nodes associated with cell τ , and ϕ_L is the shape function associated with node L , then the condition of Eq.(33) is approximated as

$$\int \hat{\rho}(\mathbf{r}) \frac{\mathbf{r} - \mathbf{R}_K}{|\mathbf{r} - \mathbf{R}_K|^{\frac{3}{2}}} d\mathbf{r} + \sum_{\tau \in \mathcal{M}} \sum_{A \in \tau} Z_A \frac{\sum_{L \in \mathcal{V}(\tau)} \mathbf{R}_L \phi_L(\mathbf{R}_A) - \mathbf{R}_K}{\left\| \sum_{L \in \mathcal{V}(\tau)} \mathbf{R}_L \phi_L(\mathbf{R}_A) - \mathbf{R}_K \right\|^{\frac{3}{2}}} = \mathbf{0} \quad (34)$$

This effectively reduces the dimension of the problem from $3M$ (the (x, y, z) coordinates of the nuclei), to $3M_{rep}$, where M_{rep} is the number of nodes in the atomic mesh (the number of *repnuclei*). The sum in Eq.(34) is most likely not going to be the simulation bottleneck (solving the electronic problem for $\hat{\rho}$ is significantly more demanding), but fast-multipole methods [22–24] can be considered to speed the summation.

Denoting by \mathbf{P}_i , $i = 1, \dots, M_{rep}$, the position of the representative nucleus n_i , the set of nonlinear equations of Eq.(34) can be grouped into a nonlinear system that is solved for the relaxed configuration of the structure.

$$\begin{aligned} \mathbf{f}_1(\mathbf{P}_1, \mathbf{P}_2, \dots, \mathbf{P}_{M_{rep}}) &= \mathbf{0} \\ \mathbf{f}_2(\mathbf{P}_1, \mathbf{P}_2, \dots, \mathbf{P}_{M_{rep}}) &= \mathbf{0} \\ \dots & \\ \mathbf{f}_{M_{rep}}(\mathbf{P}_1, \mathbf{P}_2, \dots, \mathbf{P}_{M_{rep}}) &= \mathbf{0} \end{aligned} \quad (35)$$

where \mathbf{f}_K is the left side of Eq.(34). Finding the solution of this system is done by a Newton-like method. Evaluating the Jacobian information is straightforward but not detailed here.

Finally, note that within Eq.(34) a connection is made back to Eq.(8); the position of an arbitrary nucleus A in cell τ is computed based on interpolation using the nodes $\mathcal{V}(\tau)$, one of many alternatives available (one could consider *repnuclei* from neighboring cells for instance). Effectively, this provides in Eq.(8) an expression for $\Phi(\cdot, t)$ that depends only on $A \in \mathcal{V}(\tau)$ rather than $A \in \mathcal{B}$.

PROPOSED COMPUTATIONAL SETUP

Given a nanostructure of known atomic composition (not necessarily mono-atomic or single-crystal), the goal is to determine the electron density distribution as well as the final configuration of the nanostructure, that is, the mapping Φ . Because of the assumption that the kinetic energy of the nuclei is zero, the problem corresponds to a zero Kelvin temperature scenario. A methodology that handles the nonzero temperature case is not addressed here; most likely, it would follow an approach similar to that of Car-Parrinello [25], or Payne et al. [26].

As indicated in Fig.2, the proposed solution has three principal modules: the *Preprocessing* stage, the *Electronic Problem*, and the *Ionic Problem*. *Preprocessing* is carried out once at the beginning of the simulation. A suitable chosen domain D is selected to include the nanostructure investigated. The partitioning of D into u subdomains $D_i, i = 1, \dots, u$, is done to mirror the underlying periodicity of the structure. A set of subdomains $D_{\chi(1)}$ through $D_{\chi(p)}$ is determined to constitute the reconstruction domains, and as in Fig.1, they are denoted by Y_1 through Y_p . In these p subdomains explicit electronic structure computation will be carried out accurately. A set of values of the electronic density is required at the nodes of the discretization mesh; the initial guess for the electronic density could be a uniform distribution throughout the nanostructure or, when practical, could be obtained based on a periodic boundary conditions assumption by computing it in a domain D_j and then cloning for the remaining domains D_k . *Preprocessing* concludes with the initialization of the deformation map Φ to be the identity mapping.

The *Electronic Problem* can be solved externally or internally. When it is solved externally, a specialized code such as NWChem [8] is employed to compute the electronic density in the reconstruction subdomains $Y_\alpha, \alpha \in \{1, \dots, p\}$. When the *Electronic Problem* is solved internally (only for qualitative studies, using for instance the Thomas-Fermi DFT, or OFDFT), the computation requires a mesh grid on which the integrals associated with the formulation are discretized. The algorithm uses three-dimensional interpolation to provide for the density in D_j , where $j \in \{1, \dots, u\} - \{\chi(1), \dots, \chi(p)\}$.

Independent of the type of solver invoked (external or internal), using a suitable norm the new electronic density ρ^{new}

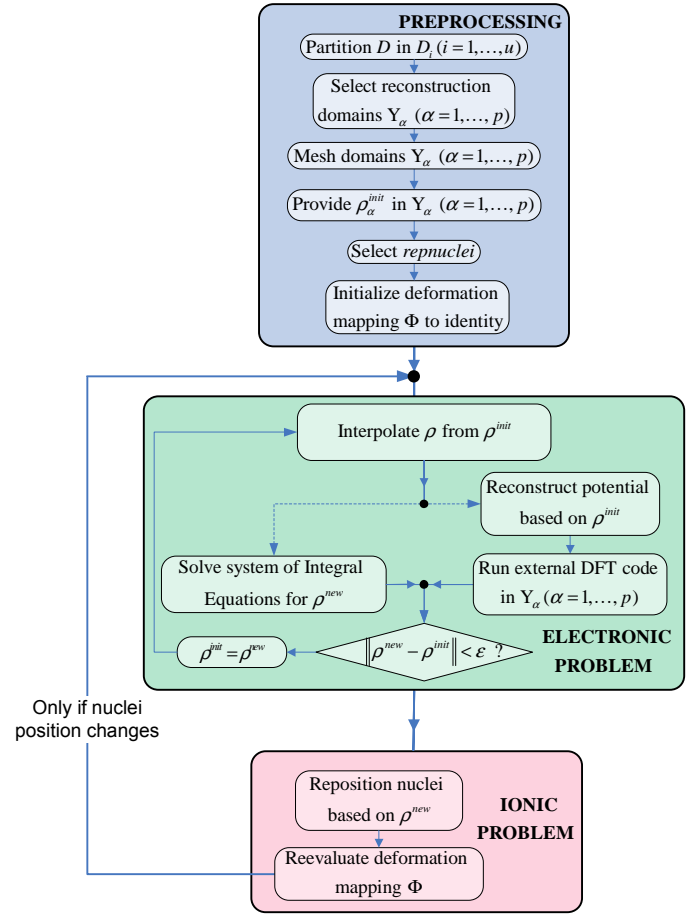


Figure 2. COMPUTATIONAL FLOW.

is compared to ρ^{init} , and the computation restarts the *Electronic Problem* after setting $\rho^{init} = \rho^{new}$, unless the corrected and initial values of the electronic density are close. This iterative process constitutes the first inner loop of the algorithm.

The *Ionic Problem* uses the newly computed electronic density to reposition the nuclei and thus alter the shape of the structure. The nonlinear system of Eq.(35) provides the position of the *repnuclei*; the other nuclei are positioned based on the quasi-continuum paradigm according to Eq.(32b). The nonlinear system in Eq.(35) is solved by an iterative method, which leads to the second inner loop that in turn has four steps:

1. Evaluate the integral of Eq.(34); when necessary, evaluate its partial with respect to \mathbf{P}_i
2. Evaluate the double sum of Eq.(34), which is based on a partitioning of the structure; when necessary, evaluate its partial with respect to the position of the representative atoms
3. Carry out a quasi-Newton step to update the positions \mathbf{P}_i of the M_{rep} representative nuclei.
4. Go back to 1 if not converged

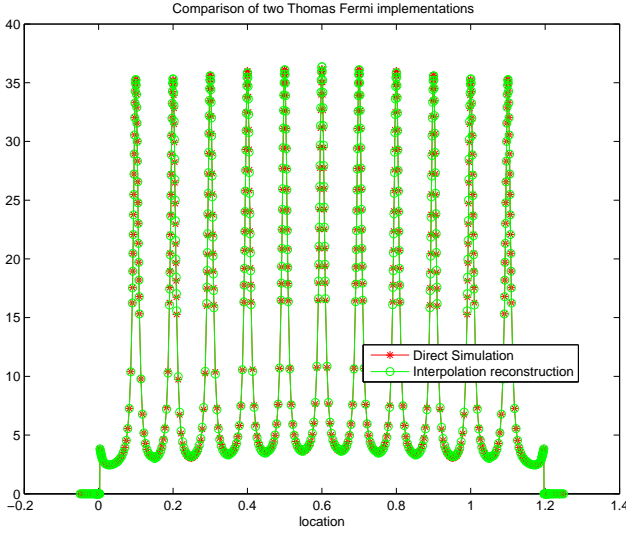


Figure 3. THOMAS-FERMI SOLUTIONS.

The precision in determining the position of the nuclei is directly influenced by the accuracy of the electronic density $\rho(\mathbf{r})$. Accurately solving the *Electronic Problem* is computationally intensive, and thus an important issue not addressed by this work is the sensitivity of the solution of the nonlinear system of Eq.(35) with respect to $\rho(\mathbf{r})$. It remains to determine whether a crude approximation of the electronic density suffices for solving the *Ionic Problem* at a satisfactory level of accuracy.

After determining the position of the nuclei, the algorithm computes the new deformation mapping Φ according to Eq.(8). If the overall change in the position of *repnuclei* at the end of the *Ionic Problem* is smaller than a threshold value, the computation stops; otherwise the new distribution of the nuclei is the input to a new *Electronic Problem* (second stage of the algorithm).

In summary, the algorithm passes through the *Preprocessing* stage once. It then solves the *Electronic Problem* (the first inner loop) and proceeds to the *Ionic Problem* (the second inner loop). The outer loop (*Electronic Problem*, followed by *Ionic Problem*) stops when there is no significant change in the position of the *repnuclei*.

NUMERICAL RESULTS

For a simple example, this section compares the numerical results obtained with the direct minimization approach of the Thomas-Fermi DFT, with the ones produced by EDR using the internal optimization approach introduced in this paper.

In the one-dimensional test, eleven equally spaced nuclei with unit charge, $Z_A = 1$, are considered; the total number of electrons is $N_e = 11$. The location of the atoms corresponds to the peaks seen in Fig.(3). A mesh is constructed that has 50 nodes per cell, with 30 of them equally spaced on an interval centered

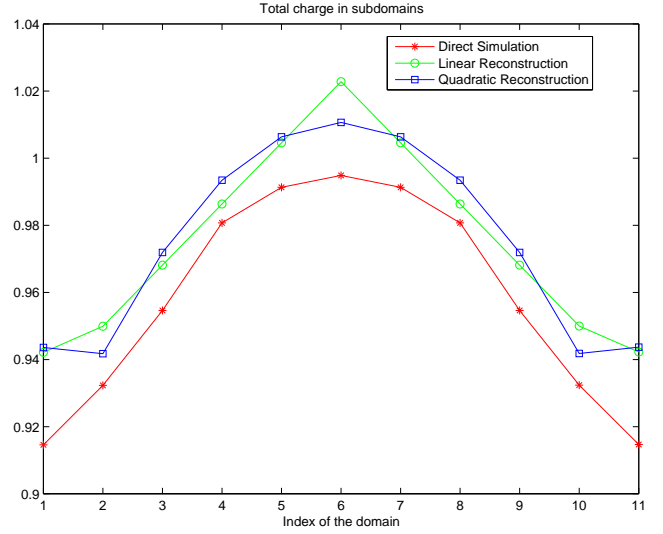


Figure 4. THOMAS-FERMI TOTAL CHARGE.

at the position of the atom and whose length is $1/5$ of the distance between two atoms. This results in 11 domains, D_1, D_2, \dots, D_{11} ; for this example, $Y_1 = D_1, Y_2 = D_2, Y_3 = D_6, Y_4 = D_{10}, Y_5 = D_{11}$. For discretization of the integral operators the trapezoidal rule was considered. Only the domains D_2, D_6, D_{10} are used for interpolation in order to avoid the boundary distortion. Either piecewise linear, or quadratic interpolations is employed for the interpolation-based approach. The parameter $\delta = 10^{-4}$ leads to a slightly different regularization from that described in the previous sections (see Eq.(19)), whereby terms of the type $1/|| \cdot ||$ are replaced with $1/|| \cdot || + \delta ||$.

The resulting electronic structure optimization problem is solved with the augmented Lagrangian software Lancelot [27], which uses an iterative method to solve the bound constrained subproblem obtained after penalization of the constraints. When using the interpolation method, the interpolation conditions of Eq.(7) are actually enforced as constraints, rather than substituting them in the functional that describes the problem (see E^{IO} of Eq.(27)). When efficiency is a concern, this substitution would be carried out and only the electronic density degrees of freedom in $D_1, D_2, D_6, D_{10}, D_{11}$ would be considered. The objective here is only to validate the interpolation-based reconstruction without regard to computational efficiency.

The solution of the direct numerical simulation and of the linear-interpolation-based optimization are presented in Fig.3. The results are almost identical, which is an indication that the interpolation approach was effective in reconstructing the solution in the “gap” domains. The results were better yet for the quadratic-interpolation-based reconstruction (not presented in the figure). Note, however, that the solutions are not identical. This can be seen by computing the total charge in the subdomains. The results for the three numerical experiments are pre-

sented in Fig.3, the right panel. This indicates that the quadratic interpolation method produced a very good fit, with a relative error that is uniformly below 2% for domains 2 through 10. The solution presents some artifacts at the very end of the domain. That in itself is not surprising, given the limitations associated with the Thomas-Fermi DFT.

CONCLUSIONS

The paper proposes a theoretical framework for nanostructure investigation. The *near-periodicity* and *small-deformation* assumptions introduced are at the center of a methodology that uses interpolation and coupled cross-domain optimization techniques to increase the size of the problems that rely on a DFT-based solution component. In this and the companion work [16], for the electron density computation (the *Electronic Problem*) formal error bounds are provided for the interpolation and cross-domain reconstruction techniques used. The EDR process can be done internally by following a nonlinear-system approach, or an optimization approach introduced here in conjunction with the Thomas-Fermi DFT; alternatively, it can be carried out externally by using dedicated third-party software such as NWChem. In either case, the EDR methodology calls for the solution of a cross-domain coupled nonlinear problem. The last step of the proposed methodology solves the *Ionic Problem*; that is, repositioning the nuclei of the structure given the electronic density in the domain. It was shown that the new ionic configuration is the solution of a nonlinear system obtained based on first-order optimality conditions. The Jacobian information for this system is readily available, and its solution does not require the explicit dependency of the kinetic and exchange-correlation energies on the electronic density. The proposed methodology is currently being implemented for three-dimensional structures. The short-term goal is to use the EDR methodology to determine the electronic structure of an inner defect in a silicon crystal. Early simulation results are promising, setting the stage for future tests with more complex structures and more sophisticated DFT approaches.

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